



Microtorus: A High Finesse Microcavity with Whispering Gallery Modes

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Microsphere -- a low-loss photon trap, efficient optical cavity

Whispering-gallery modes - closed circular waves under total internal reflection

(Term by J.W.S.Rayleigh, analogy to acoustic modes in the gallery of St Paul cathedral)

(MUST BE) Sustained in any axisymmetric dielectric body with $R \gg \lambda$

low material loss (transparent material, e.g fiber grade silica)

low bending loss ($R \gg \lambda$)

low scattering loss (TIR always under grazing incidence

+ molecular-size surface roughness)

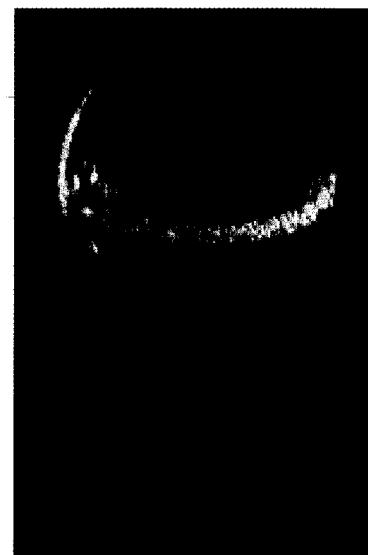
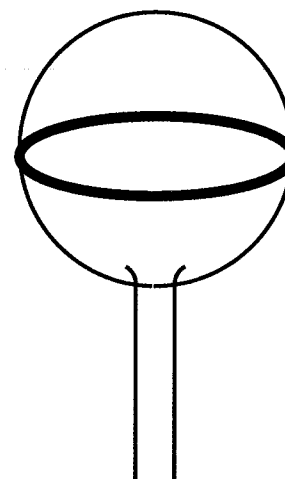
Quality-factor $Q = \lambda / \Delta\lambda_{\text{RES}}$ - up to $\sim 10^{10}$

Photon lifetime $\tau = \lambda Q / 2\pi c$ - up to $\sim 3\mu\text{s}$

(cavity ringdown time)

visible and near-infrared band: *Opt.Lett. 21, p.453 (1996)*

Opt.Lett. 23, p.247 (1998)



Visualization of WG mode field by residual scattering in silica microsphere, *V.S.Ilchenko et al, Opt.Comm. 113, p.133(1994)*



Why spheres?

low material loss (transparent material)

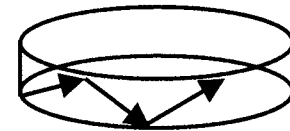
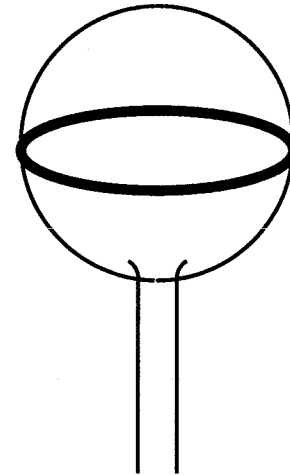
low bending loss (high-contrast boundary)

LOW SCATTERING LOSS (TIR always under grazing incidence) $\Theta \rightarrow \pi/2$; compare to disks/ rings:

$$\frac{I_R}{I_I} = e^{-\left(\frac{4\pi\sigma}{\lambda} \cos\Theta\right)^2} \quad (\text{J.W.S.Rayleigh})$$

EVEN WITH MOLECULAR ROUGHNESS σ , ONLY CURVATURE CONFINEMENT ALLOWS **Q** LIMITED BY MATERIAL ATTENUATION:

$10^8 \dots 10^{10}$ in spheres vs. $10^3 \dots 10^5$ in microrings !!



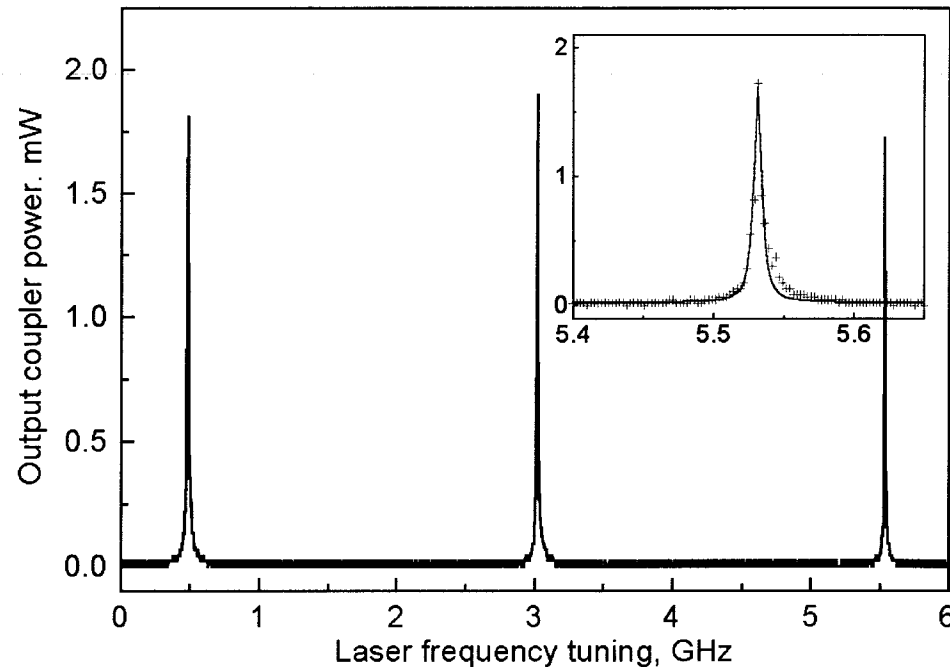
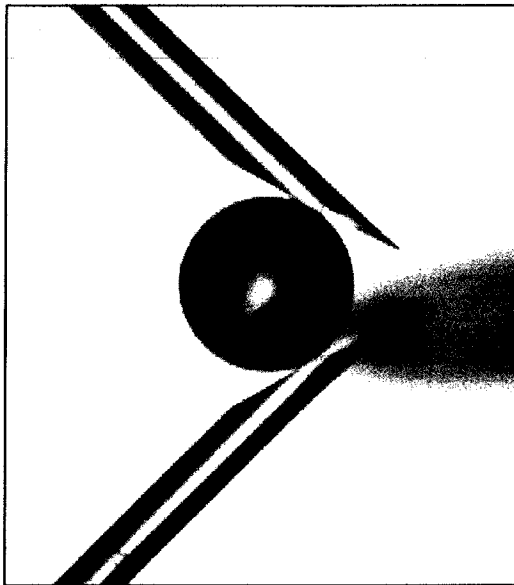
Drawback: “too many modes” compared to planar rings!



Spectrum of microspheres: Families of non-degenerate $TE(TM)_{lmq}$ modes.

“Small” FSR $\nu_{lmq} - \nu_{l,m-1,q} \sim \nu \frac{\varepsilon^2}{2l}$ - few GHz with typical $\varepsilon^2 \sim (1-3) \times 10^{-2}$.

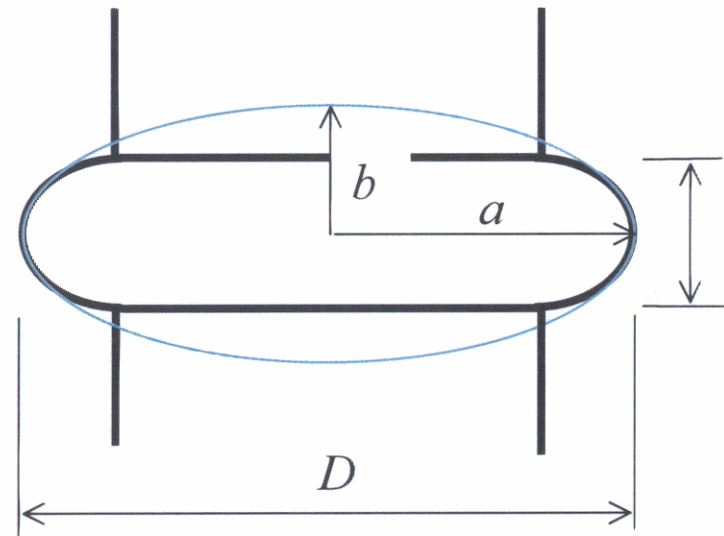
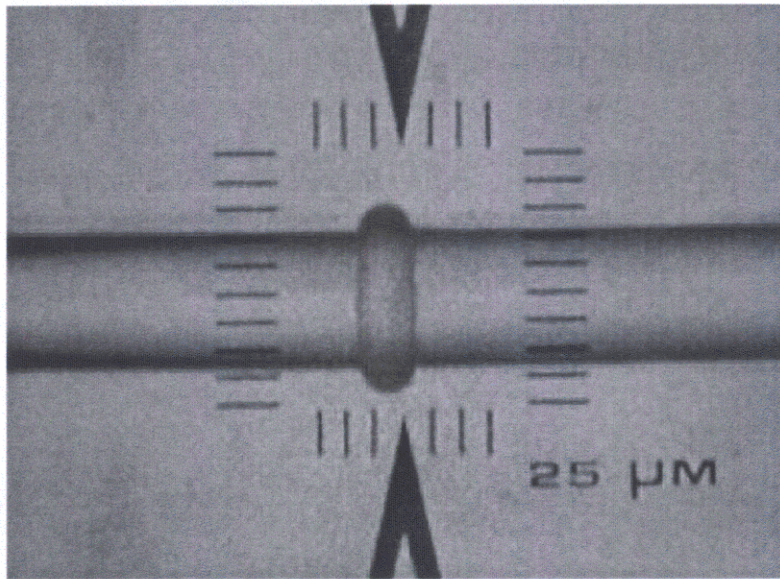
“Big” FSR $\nu_{lmq} - \nu_{l-1,mq} = \frac{c}{2\pi na} (t_{lq} - t_{l-1,q}) \sim \nu / l$ - few hundred GHz (few nm)



Input power 7.5...8.3mW; maximum transmission at resonance $\sim 23.5\%$ (fiber-to-fiber loss 6.3dB); $Q_{load} > 3 \times 10^7$ at 1550nm; sphere diameter 405 μ m. Unloaded $Q_o \approx 1.2 \times 10^8$ (*Opt.Lett.* 24, 723 (1999))



Novel geometry: a highly oblate spheroid, or microtorus



Near the symmetry plane (at the location of WG modes), toroidal surface of outer diameter D and cross-section diameter d coincides with that of the osculating oblate spheroid with large semiaxis $a = D / 2$ and small semiaxis $b = \frac{1}{2} \sqrt{Dd}$



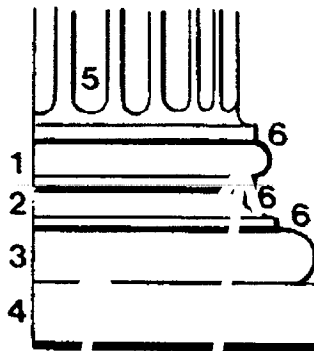
Names

[F, seesaw]: an apparatus or structure
end is counterbalanced by the other on
or by weights

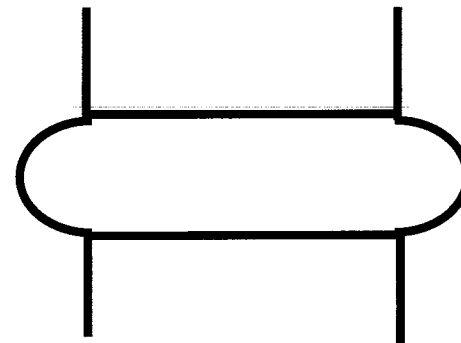
\'bā-səz\ [ME, fr.
sp, base, fr. *bainein*

a : the bottom of
support : FOUNDA-
t of a wall, pier, or
parate architectural
part of a complete
side or face of a
ich an altitude can
on which the figure
a bodily organ by
other more central

2 a : a main in-
latex ~> **b** : a
ingredient (as of a
ental part of some-
the lower part of a
oint or line from which a start is made
b : a line in a survey which serves as



base of a column:
1 upper torus, 2
scotia, 3 lower
torus, 4 plinth, 5
shaft, 6 fillets



(Webster's New College dictionary, G & C Merriam Co. Springfield, Mass., 1975, p.92)



Calculation of the spectrum of the dielectric spheroid

is not a trivial problem, even numerically. In “quasiclassical” approximation with assumptions:

1) a WG mode is a closed circular beam supported by TIR, 2) optical field tunnels outside at the depth $1/k\sqrt{n^2-1}$, and 3) the tangential component of E (TE -mode), or normal of D (TM -mode) is continuous at the boundary. Eigenfrequencies of high-order WG modes ($l \gg 1; l \approx m$) in dielectric sphere can be approximated via solutions of scalar wave equation with zero boundary conditions, because most of the energy is concentrated in one component of the field (E_θ for TE -mode and E_r for TM -mode).

Based on above considerations, let us estimate WG mode eigenfrequencies in oblate spheroids of large semiaxis a , small semiaxis b , and eccentricity $\varepsilon = \sqrt{1-b^2/a^2}$. Since WG modes are localized the “equatorial” plane, we shall approximate the radial distribution by cylindrical Bessel function $J_m(n\tilde{k}_{mq}r)$ with $n\tilde{k}_{mq}a = na\sqrt{k_{lmq}^2 - k_\perp^2} \approx T_{mq}$, where $J_m(T_{mq}) = 0$ and k_\perp is the wavenumber for quasiclassical solution for angular spheroidal functions. For our purposes a rough approximation is enough: $k_\perp^2 \approx \frac{2(l-m)+1}{a^2\sqrt{1-\varepsilon^2}} m$;

more rigorous consideration can follow the approach given in [I.V.Komarov, L.I.Ponomarev, S.Yu.Slavyanov, *Spheroidal and Coulomb Spheroidal Functions*, Moscow, Nauka (1976) (in Russian)]. Taking into account that $T_{mq} \approx t_{lq} - (l-m+1/2)$, we finally obtain the following approximation:

$$nk_{lmq}a - \frac{\chi}{\sqrt{n^2-1}} \approx t_{lq} + \frac{2(l-m)+1}{2} \left(\frac{1}{\sqrt{1-\varepsilon^2}} - 1 \right)$$

t_{lq} -- q -th zero of the spherical Bessel function of the order l ; $\chi = n$ for TE -mode, $\chi = 1/n$ for TM -mode.

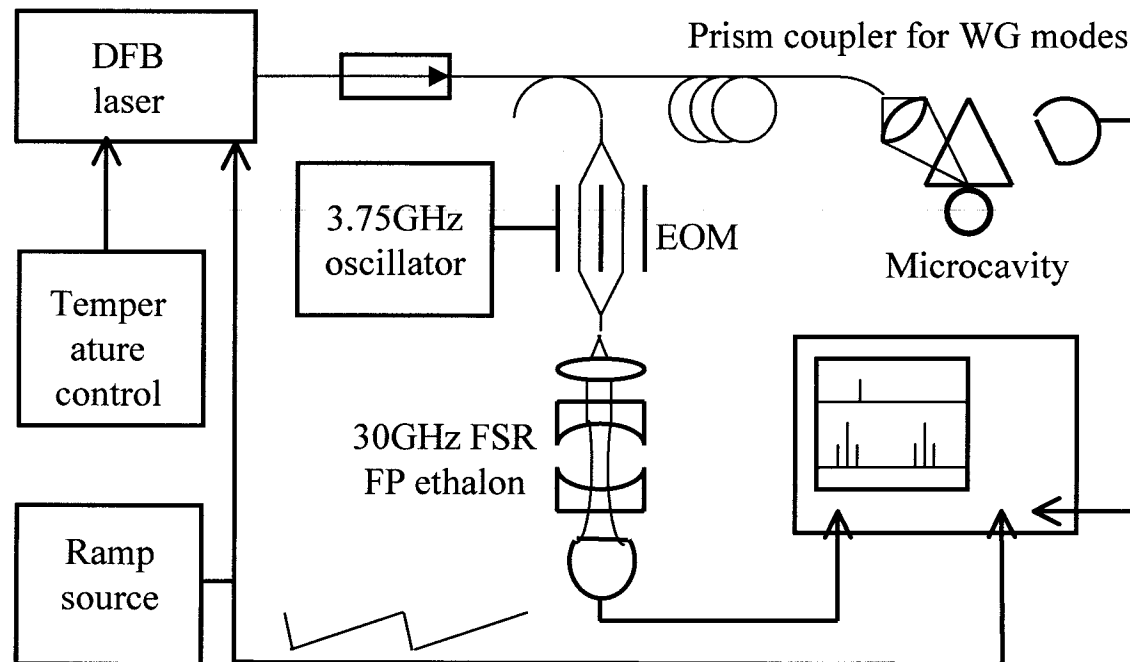
1. For small eccentricities the model gives identical prediction with perturbation theory (H.M.Lai, P.T.Leung, K.Young, P.W.Barber, S.C.Hill, *Phys. Rev. A* 41, 5187-5198 (1990))

2. Discrepancy with numerical calculations is <5% in prediction of “small” FSR and <0.1% of absolute frequencies, even with $\varepsilon^2 \sim 0.8$, even small $l = 100$



Schematic of the experimental setup

to obtain wide range ($\sim 900\text{GHz}$, or 7.2nm) high-resolution spectra of WG modes in microcavity

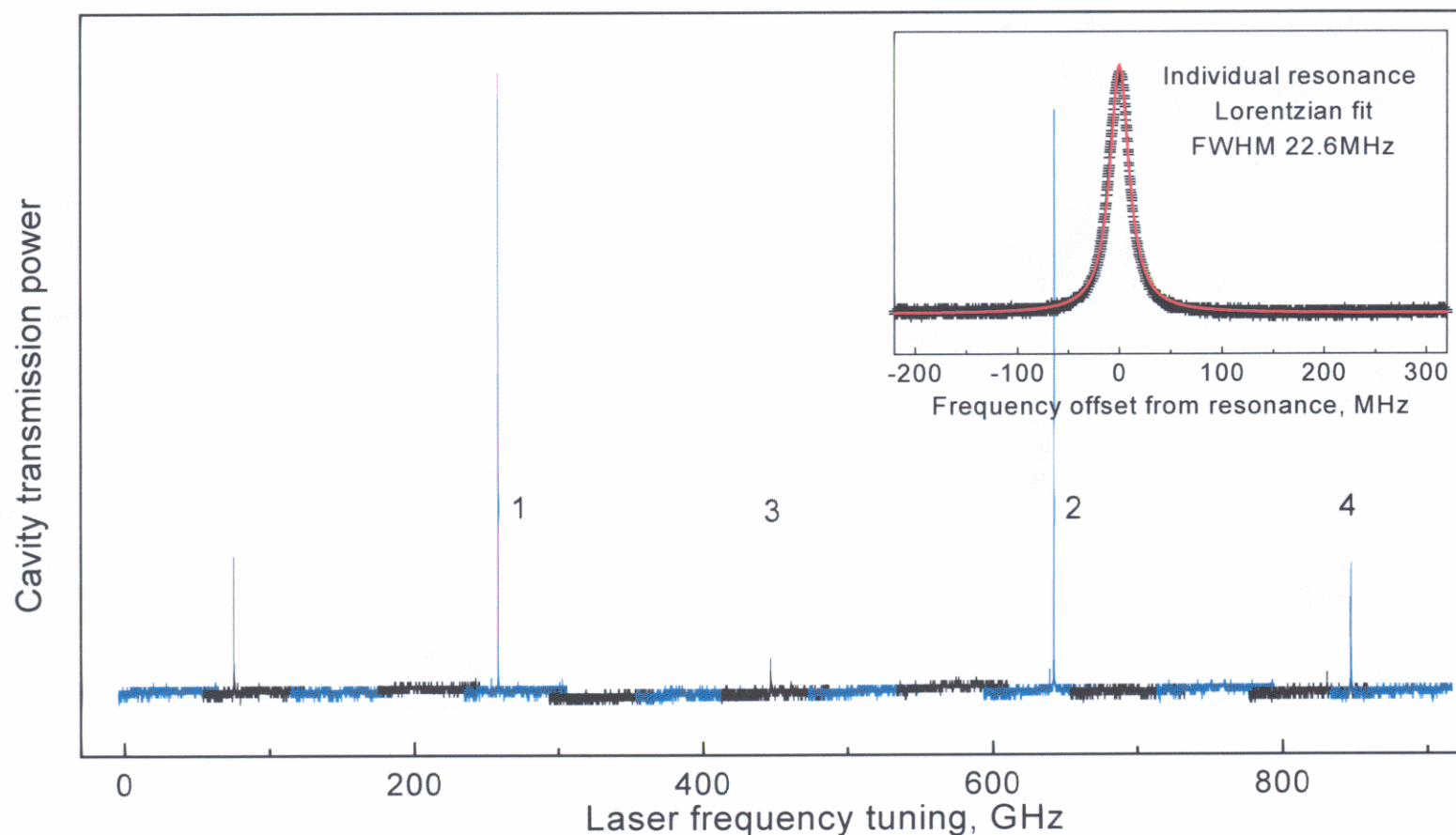


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Spectrum of TE whispering-gallery modes in spheroidal dielectric microcavity

$D = 2a = 165\mu\text{m}$; $d = 42\mu\text{m}$; $2b = 83\mu\text{m}$. Free spectral range (between largest peaks 1 and 2) 383.7GHz (3.06nm) near central wavelength 1550nm. Individual resonance bandwidth 23MHz (loaded $Q = 8.5 \times 10^6$).

$$\text{Finesse } F = 1.7 \times 10^4$$





CONCLUSIONS

1. It seems indeed we can combine small size, ultra-high-Q with “nice” FP-like spectrum: true finesse $10^4 \dots 10^6$ becomes available in microcavities as opposed to supermirror FPs
2. Host of potential applications is significant